

# Light-induced thermoelastic sensor for ppb-level H<sub>2</sub>S detection in a SF<sub>6</sub> gas matrices exploiting a mini-multi-pass cell and quartz tuning fork photodetector

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## ABSTRACT

We present an optical sensor based on light-induced thermoelastic spectroscopy for the detection of hydrogen sulfide (H<sub>2</sub>S) in sulfur hexafluoride (SF<sub>6</sub>). The sensor incorporates a compact multi-pass cell measuring 6 cm × 4 cm × 4 cm and utilizes a quartz tuning fork (QTF) photodetector. A 1.58 μm near-infrared distributed feedback (DFB) laser with an optical power of 30 mW serves as the excitation source. The sensor achieved a minimum detection limit (MDL) of ~300 ppb at an integration time of 300 ms, corresponding to a normalized noise equivalent absorption coefficient (NNEA) of  $3.96 \times 10^{-9} \text{ W} \cdot \text{cm}^{-1} \cdot \text{Hz}^{-1/2}$ . By extending the integration time to 100 s, the MDL can be reduced to ~25 ppb. The sensor exhibits a response time of ~1 min for a gas flow rate of 70 sccm.

## 1. Introduction

Sulfur hexafluoride (SF<sub>6</sub>) is a colorless, odorless, non-toxic, non-flammable chemically inert gas. It possesses excellent dielectric properties, making it an exceptional arc-quenching medium. When compared to air, SF<sub>6</sub> exhibits a significantly higher arc-extinguishing ability, being 100 times more effective. Moreover, SF<sub>6</sub> demonstrates superior insulation capability, which is 2.5 times better than that of air, along with a density that is 5.6 times greater. Consequently, SF<sub>6</sub> has emerged as a replacement for oil and compressed air as an insulation medium in power grid systems. Applications of SF<sub>6</sub> include gas-insulated switchgears, transmission pipes, gas circuit breakers, and transformers [1–3]. In power grid, gas-insulated switchgears (GIS) play a crucial role in facilitating high-voltage switch on/off operations exceeding 1 MV [4]. Nevertheless, over prolonged periods of operation, corona, partial discharge, and spark discharge may transpire within the GIS. SF<sub>6</sub> is typically stable under normal conditions. However, in the presence of metals within an overheated environment, discharge events can lead to the generation of various decomposition compounds [5,6], such as SF<sub>4</sub>, SF<sub>3</sub>, SF<sub>2</sub>, and S<sub>2</sub>F<sub>10</sub>. These low-fluorine sulfides have the potential to

react with minute amounts of air and moisture present within the equipment, resulting in the formation of toxic and corrosive compounds, such as H<sub>2</sub>S, HF, SOF<sub>2</sub>, SO<sub>2</sub>F<sub>2</sub>, SOF<sub>4</sub>, SO<sub>2</sub> and CO [7–10]. Monitoring the concentrations and formation rates of these decomposition products can provide insights into the type and progression of internal faults within the equipment [11]. The detection of hydrogen sulfide (H<sub>2</sub>S) traces holds particular significance in assessing the discharge energy of a fault and determining whether the fault involves solid insulation materials [12]. Therefore, achieving detection sensitivity in the sub ppm-range for H<sub>2</sub>S in SF<sub>6</sub> becomes crucial for ensuring the security maintenance of the power grid [13].

In recent years, various techniques have been employed for the detection of decomposition compounds formed within the SF<sub>6</sub> matrix. These techniques include tube detection, metal-oxide semiconductor sensing, gas chromatography [14–16] and, more recently, in situ and real-time optical sensors [17,18]. In 2017, the first demonstration of a photoacoustic spectroscopy (PAS) sensor for H<sub>2</sub>S detection in SF<sub>6</sub> was conducted [13]. PAS employs a microphone to detect the faint sound waves produced as a result of molecule absorption through the photoacoustic effect. However, the use of a highly sensitive microphone with a

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wide acoustic detection bandwidth presents a challenge when installing a PAS sensor in a noisy environment, as any external acoustic source can contribute to the overall signal. As a variant of PAS, quartz enhanced photoacoustic spectroscopy (QEPAS) is particularly well-suited for in situ and real-time measurement applications owing to its robust immunity to environmental noise [19]. This is achieved by utilizing a quartz tuning fork (QTF) in place of the traditional microphone for sound wave detection. The QTF operates at resonant frequencies with a narrow bandwidth ( $<1$  Hz), effectively filtering out any noise contributions outside of this range and ensuring high immunity to environmental noise. In QEPAS, the quartz tuning fork (QTF) is completely immersed in the fluid containing the target molecules that need to be detected. However, this immersion poses limitations on the application of QEPAS techniques for the detection of decomposition compounds within the  $\text{SF}_6$  matrix for two primary reasons. First, the heavy molecular weight of  $\text{SF}_6$  molecules results in significant damping of the vibration amplitudes of the QTF within the  $\text{SF}_6$ -based matrix, as compared to the behavior observed in standard air [20]. Second, decomposition substances such as  $\text{SO}_2$  and  $\text{H}_2\text{S}$  in  $\text{SF}_6$  exhibit high corrosiveness, particularly towards the metal electrodes coated on the surfaces of the QTF. Consequently, this presents a constraint on the long-term operation of the QEPAS system. [21,22].

Recently, QTFs have been proposed as a sensitive and cost-effective alternative to infrared photodetectors in optical sensors that utilize light-induced thermoelastic spectroscopy (LITES) [23–27]. The LITES setup architecture closely resembles that of tunable diode laser spectroscopy combined with the wavelength modulation technique. In this setup, an intensity-modulated laser beam traverses a gas cell, and subsequently, the light emerging from the gas cell is focused onto one of the prongs of the QTF, inducing localized modulated heating. This thermal energy generates a local strain field, which in turn gives rise to a piezoelectric charge distribution. Similar to QEPAS, this charge distribution can be collected using the metal contact patterns deposited on the surface of the QTF. Thus, both QEPAS and LITES utilize the properties of the QTF, such as the high resonance Q-factor, narrow frequency bandwidth, wavelength independence, and the possibility of implementing a wavelength modulation approach with second harmonic detection to achieve a noise-free background [24]. The main difference between QEPAS and LITES lies in the physical phenomenon that causes the oscillation of the QTF prongs and consequently the production of an electric charge distribution. In contrast to QEPAS, LITES offers the distinct advantage of having the QTF located outside the gas cell, enabling non-contact measurements [25–32].

Here we report on the implementation of a LITES sensor for detecting  $\text{H}_2\text{S}$  in  $\text{SF}_6$  in order to monitor gas-insulated switchgears. The vibrational characteristics of a standard 32 kHz-QTF are analyzed in different environments ( $\text{SF}_6$  and  $\text{N}_2$ ), validating the benefits of employing LITES technology compared to QEPAS for  $\text{SF}_6$  decomposition monitoring. The operating parameters of the developed LITES sensor are optimized, and the obtained results are thoroughly discussed.

## 2. Vibration analysis of QTF in $\text{N}_2$ and $\text{SF}_6$ environments

The QTF is a mechanically forked structure composed of a quartz crystal with piezoelectric properties. When it absorbs energy from sound waves or a modulated laser beam that resonates with the frequency of the QTF, the prongs of the fork initiate vibration, resulting in the generation of a strain field. Utilizing the piezoelectric effect, a charge distribution is produced and collected through the electrode pattern distributed on the surface of the prongs. This electrical signal can be demodulated using a lock-in amplifier, allowing for the extraction of a signal that is directly proportional to the intensities of the sound waves (QEPAS) or the absorbed laser beam (LITES).

The resonant frequency of the QTF is determined by the prongs geometry, and the resonant frequency can be estimated by the following equation [32]:

$$f = \frac{1}{2\pi} \sqrt{\frac{k}{m_{\text{eff}}}} \approx 1.015 \frac{w}{2\pi L^2} \sqrt{\frac{E}{\rho}} \quad (1)$$

where  $k = Etw^3/(4L^3)$  is stiffness,  $m_{\text{eff}} = 0.2427\rho Lwt$  is effective mass,  $E$  and  $\rho$  are Young's modulus and density of quartz, respectively,  $L$ ,  $w$ ,  $t$  are the prongs length, width, and thickness, respectively.  $L$ ,  $w$ ,  $t$  for a commercial standard 32-kHz QTF were measured to be 3.8 mm, 1.4 mm and 0.3 mm, respectively. For certain gases with low relaxation rates such as methane, carbon monoxide, and nitric oxide, the frequency of a standard 32-kHz QTF may not satisfy the requirement  $\tau \ll 1/f$  ( $\tau$  is molecular vibration-rotational relaxation rate) needed for effective generation of electrical signals. However, when measuring  $\text{H}_2\text{S}$  concentration, which is a fast-relaxing molecule, the 32-kHz QTF can be operated successfully [33]. In addition to its small size, the standard 32-kHz QTF offers several advantages including affordability, low signal loss, wide dynamic range, robustness against environmental noise, and high temperature resistance [34–37]. Nonetheless, its applicability for detecting  $\text{SF}_6$  decomposition may be limited due to the high density of  $\text{SF}_6$  in comparison to standard air or  $\text{N}_2$ .

When a QTF undergoes small amplitude harmonic oscillations while immersed in a fluid medium, it induces movement in the surrounding fluid. This phenomenon leads to energy dissipation through acoustic losses and introduces additional inertia. As a result, the reaction force comprises a resistive component responsible for energy dissipation and a reactive component that contributes to the increased inertia experienced by the vibrating QTF. Consequently, the vibration of a QTF's prong can be described by a fourth-order differential equation under these conditions. [32]:

$$EI \frac{\partial^4 y}{\partial x^4}(x, t) + C_d \frac{\partial y}{\partial t}(x, t) + (\rho A + u) \frac{\partial^2 y}{\partial t^2}(x, t) = 0 \quad (2)$$

where  $\rho$  represents the density of the quartz material,  $E$  stands for the Young's modulus of the material,  $t$  denotes time,  $C_d$  is the damping parameter,  $u$  is the added mass per unit length,  $A$  is defined as the cross product of the crystal thickness  $w$  and the prong width  $T$ , and the directions  $x$  and  $y$  refer to the plane of the QTF. For small damping, it has been demonstrated that the presence of an added mass causes a red shift  $\Delta f$  in the resonance frequency of the QTF [34], in comparison to its resonance frequency in a vacuum:

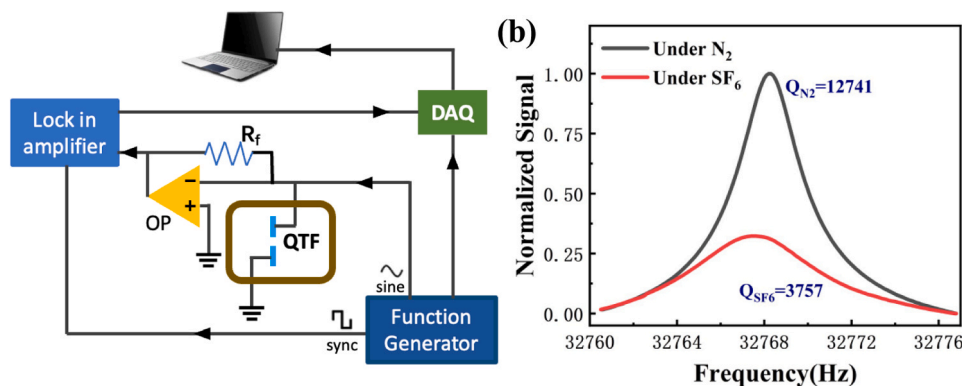
$$\Delta f = -\frac{1}{2} \frac{u}{\rho A} \quad (3)$$

The added mass is directly proportional to the density  $\rho_0$  of the gas medium and, consequently, its pressure. Thereby when operating the QTF in  $\text{SF}_6$  compared to  $\text{N}_2$  or standard air, larger red shifts are expected.

The Q-factor of a resonance mode is a measure of the energy loss experienced by the prongs while vibrating. To model the damping effect caused by air, an analytical expression derived by Hosaka et al. [38] can be utilized, allowing for the formulation of the contribution of gas damping to the Q-factor:

$$Q_{\text{gas}} \approx \frac{4\rho T w^2 f}{3\mu w + \frac{3}{4}w^2 \sqrt{4\pi\rho_0\mu f}} \quad (4)$$

where  $\mu$  is the gas viscosity. Gas damping has been identified as the primary mechanism responsible for energy loss, both in the case of  $\text{N}_2$  and  $\text{SF}_6$  [20]. The density of  $\text{N}_2$  and  $\text{SF}_6$  are  $1.16 \text{ kg/m}^3$  and  $6.52 \text{ kg/m}^3$  [15], thereby, a lower QTF Q-factor is expected when it is placed in a  $\text{SF}_6$  matrix with respect to  $\text{N}_2$  or standard air. To evaluate the decrease in the Q-factor of a 32 kHz QTF when operated in  $\text{SF}_6$  compared to  $\text{N}_2$ , the experimental arrangement illustrated in Fig. 1(a) was utilized. A high-resolution function generator was utilized to generate a sine voltage spanning the frequency range of 32760–32770 Hz for exciting the QTF. The current output from QTF was converted to voltage through an operational amplifier and subsequently demodulated using a lock-in



**Fig. 1.** (a) schematic of the setup used for QTF frequency scanning. OP: operational amplifier, DAQ: data acquisition, feedback resistor:  $R_f = 10 \text{ M}\Omega$ . (b) Frequency response curves of the 32 kHz-QTF at atmospheric pressure in  $\text{N}_2$  and  $\text{SF}_6$ .

amplifier (Stanford Research Model SR830). The lock-in amplifier was configured with parameters including a filter slope of 12 dB/oct and a time constant of 100 ms, resulting in a detection bandwidth of 1.6675 Hz. During the measurements, an integration time of 300 ms was employed. The resonance curves of a 32 kHz-QTF immersed in  $\text{SF}_6$  or  $\text{N}_2$  at atmospheric pressure are displayed in Fig. 1(b). The QTF signal at the resonance frequency in  $\text{SF}_6$  is 3.2 times smaller than that in  $\text{N}_2$ , the quality factor drops from 12741 to 3757, and, as expected, the resonance frequency red shifts.

### 3. LITES $\text{H}_2\text{S}$ detector in $\text{SF}_6$ setup

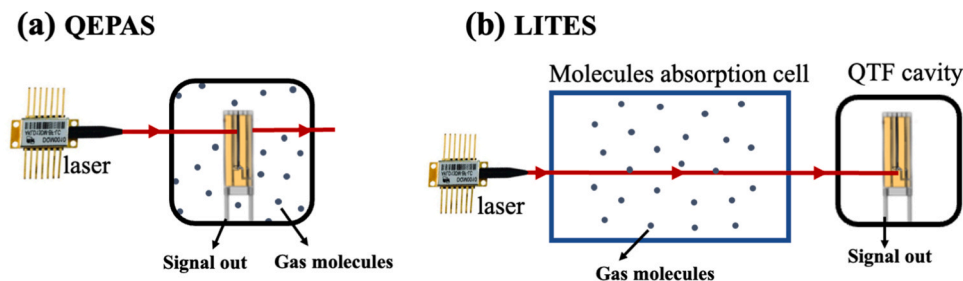
In contrast to QEPAS, the LITES technique offers a non-contact sensing approach with comparable sensitivity levels [39–41]. A scheme of the configuration of a standard QEPAS and LITES technique is reported in Fig. 2(a) and (b), respectively.

In the LITES technique, the QTF operates as an infrared photodetector, offering a cost-effective alternative to expensive photodetectors while eliminating the need for thermoelectric cooling systems [40]. A schematic of the complete architecture of the LITES sensor for  $\text{H}_2\text{S}$  detection in a  $\text{SF}_6$  matrix is depicted in Fig. 3. For the excitation source, a low-cost distributed feedback (DFB) pigtailed laser (FITEL, model FRL 15DCWD-A82), emitting within the range 1563–1583 nm, was employed. The laser is securely mounted on the driver board and controlled by a computer to maintain a stable temperature and set the laser current. The DFB laser targets a specific  $\text{H}_2\text{S}$  absorption peak occurring at  $6320.55 \text{ cm}^{-1}$ . Despite the low line strength of the targeted  $\text{H}_2\text{S}$  absorption feature, which measures  $1.1 \times 10^{-22} \text{ cm}^2/\text{mol}$ , no fiber amplifier was utilized in contrast to previous experiments [15,33]. This simplifies the setup architecture, reduces the overall cost, and minimizes the size of the sensor system. A function generator is used to modulate the laser current at the half of the QTF's fundamental frequency. The function generator is additionally responsible for supplying a voltage ramp to the laser, allowing for a gradual scanning of the laser

wavelength across the desired absorption feature. The fiber output beam is then coupled with a non-commercial mini multi-pass cell (MMPC) that contains the target gas sample. The MMPC consists of two spherical mirrors and is designed to generate a spot pattern with 7-nonintersecting-circle [18]. Within the MMPC, the laser beam undergoes over a hundred reflections between two high-reflection mirrors. Subsequently, it exits through an aperture in the output mirror and illuminates the QTF near to one of the prongs base, precisely at 3.8 mm from the QTF's prong top, where the strongest LITES signal is predicted [30,42]. The QTF signal is then directed to a low-impedance amplifier with a feedback resistor of  $10 \text{ M}\Omega$ . Its output signal is subsequently fed into a lock-in amplifier (Stanford Research Systems, Model SR830) for the demodulation at the QTF resonance frequency, utilizing wavelength modulation and second harmonic detection. The lock-in amplifier parameters employed include a filter slope of 12 dB/oct, a time constant of 100 ms, and an integration time of 300 ms, which are consistent with previous experiments.

The MMPC is positioned within a chamber that equipped with a pre-aligned fiber collimator and a fiber connector (see Fig. 3). These components facilitate the coupling of the pigtailed laser on one side and enable the laser beam to exit through a quartz window with a diameter of 1 cm on the opposite side. To regulate the pressure inside the chamber and the gas flow, a pressure meter and a flux meter are employed. To generate various concentrations of  $\text{H}_2\text{S}$  in  $\text{SF}_6$ , a gas mixing system is utilized, starting from two certified standard cylinders containing pure  $\text{SF}_6$  and  $\text{SF}_6$  with a concentration of 100 ppm  $\text{H}_2\text{S}$ . These gas mixtures are then introduced into the MMPC chamber for sensor analysis, testing, and calibration purposes. Within the experimental setup, the QTF is housed in a small cell with a volume of  $6 \text{ cm}^3$ , featuring an optical window. This cell is connected to a diaphragm pump and a pressure meter via a needle valve, allowing for the adjustment and monitoring of the QTF's operating pressure.

Compared to the traditional LITES technology that utilizes a gas cell as absorption path [23], the MMPC enables the achievement of a



**Fig. 2.** (a) Sketch of a standard QEPAS technique. The QTF is immersed in the gas cell. (b) Sketch of a standard LITES technique. The QTF operates as an infrared photodetector outside of the gas cell.

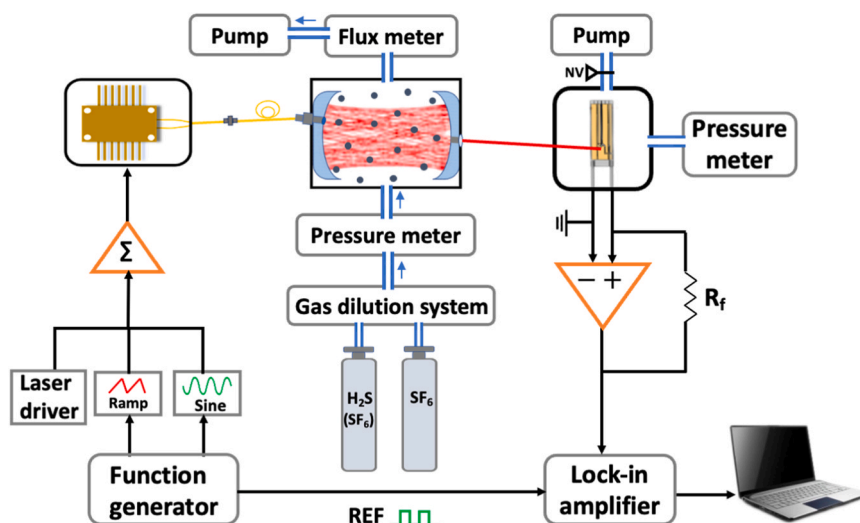


Fig. 3. Schematic of the LITES sensor for  $\text{H}_2\text{S}$  detection in the  $\text{SF}_6$  matrix. NV: needle valve; Rf: feedback resistor.

significantly longer pathlength through multiple reflections of laser beam from the mirrors. The MMPC consists of two 1-inch  $\text{CaF}_2$  spherical mirrors with a reflectivity of 99% and a radius of curvature of 2.5 cm. It is constructed with a cage structure measuring  $6\text{ cm} \times 4\text{ cm} \times 4\text{ cm}$ , providing an internal sample volume of  $67\text{ cm}^3$ . All brackets of the MMPC are made of carbon fiber material, which offers advantages such as high strain resistance, lightweight properties, and the ability to operate in extreme environments and ultra-high temperatures. The fiber collimator, equipped with a G-lens, is integrated on the left mirror frame (see Fig. 3). By utilizing the FC/APC interface, the fiber laser can be connected to the MMPC at an angle of  $8^\circ$  relative to the optical axis. Subsequently, the laser beam undergoes multiple reflections back and forth between the mirrors for a total of 107 times before eventually exiting through a small hole on the right mirror. A picture of the multi-pass cell light point/pattern is reported in Ref. [18]. Given the distance between the mirrors is  $\sim 3.9\text{ cm}$ , the total optical path length amounts to  $\sim 4.2\text{ m}$ .

#### 4. System performance assessment

When the temperature of the laser driver is set to  $31.3^\circ\text{C}$  and the output current to 148 mA, the wavenumber emission is  $6320.55\text{ cm}^{-1}$  with an optical power of 30 mW. The function generator was utilized to enable periodic scanning of the  $6320\text{ cm}^{-1}$  -  $6321\text{ cm}^{-1}$  range by the laser, targeting the selected  $\text{H}_2\text{S}$  absorption line [43]. Since the performance of the LITES sensor operating in wavelength modulation depends on the current modulation amplitude, this parameter was preliminarily optimized. The peak-to-peak value of the current modulation maximizing the LITES signal was determined to be 19 mA. In the next step, the photodetection performance of the QTF was evaluated as a function of the operating pressure. As shown in Fig. 4, the WM LITES signal was measured for a gas mixture of 100 ppm  $\text{H}_2\text{S}/\text{SF}_6$  at two different representative QTF pressures: 60 Torr and 5 Torr. The gas flow rate within the MMPC remained constant at 70 standard cubic centimeters per minute (sccm) for both measurements, while the pressure within the MMPC was maintained at atmospheric pressure (760 torr). It was observed that as the QTF pressure decreased, the LITES signal increased, reaching its highest value when operating at the lowest achievable pressure of 5 Torr. Under this condition, the LITES signal nearly doubled compared to the signal measured at 60 Torr. This behavior can be attributed to the fact that the LITES signal is proportional to the accumulation time of the oscillator, which is in turn related to the quality factor (Q-factor) of the QTF. The Q-factor is inversely proportional to the operating pressure. Based on this observation, all subsequent

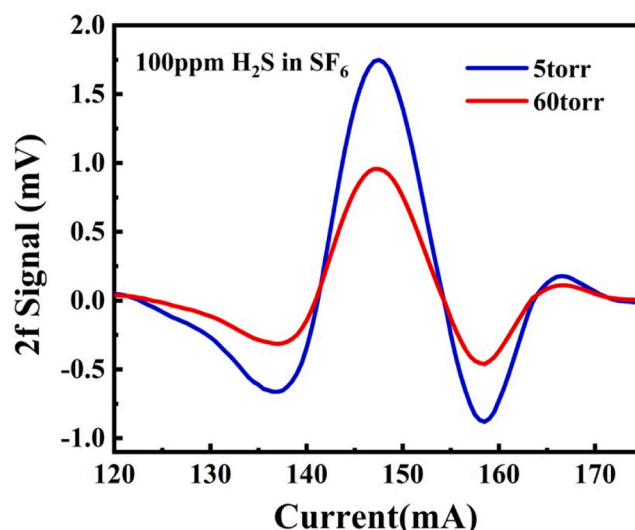


Fig. 4. Normalized WM LITES spectra measured for a mixture of 100 ppm  $\text{H}_2\text{S}$  in  $\text{SF}_6$  at two different operating pressures of the QTF-photodetector: 5 Torr and 60 Torr.

experiments were conducted with the QTF operating at 5 Torr. The performance of the  $\text{H}_2\text{S}$  in  $\text{SF}_6$  sensor system was then investigated.

Starting with the certified concentration of 100 ppm  $\text{H}_2\text{S}$  in  $\text{SF}_6$ , the LITES signal was measured at various  $\text{H}_2\text{S}$  concentrations generated using a gas dilution system by keeping the flow rate always at 70 sccm. The spectral scans of the  $\text{H}_2\text{S}$  absorption line at different  $\text{H}_2\text{S}$  concentrations in  $\text{SF}_6$  are reported in Fig. 5(a). For each spectral scan, the peak values were extracted and plotted as a function of  $\text{H}_2\text{S}$  concentrations in Fig. 5(b). The data were fitted with a linear function, yielding a slope of  $21\text{ }\mu\text{V/ppm}$  and an  $R^2$  value of 0.9996, indicating a highly linear sensor response to  $\text{H}_2\text{S}$  concentrations. The inset in Fig. 5(a) displays the background noise level of the system, measured in a spectral region with no  $\text{H}_2\text{S}$  absorption features. The  $1\text{-}\sigma$  noise level was  $\sim 6\text{ }\mu\text{V}$  for an integration time of 300 ms and the minimum detection limit (MDL) of the sensor was  $\sim 300\text{ ppb}$ . The normalized noise equivalent absorption (NNEA) coefficient was determined to be  $3.96 \times 10^{-9}\text{ W}\cdot\text{cm}^{-1}\cdot\text{Hz}^{-1/2}$ .

To estimate the  $1\text{-}\sigma$ -noise level as a function of the integration time, an Allan-Werle deviation analysis was conducted. During the analysis, the laser was locked onto the selected  $\text{H}_2\text{S}$  absorption peak while pure



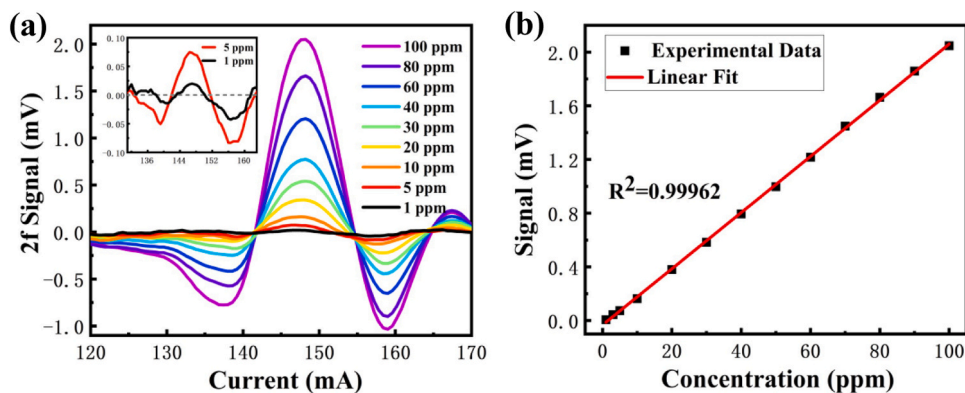


Fig. 5. (a) WM LITES spectra at different  $\text{H}_2\text{S}$  concentrations in  $\text{SF}_6$ , as reported. The laser current was swept from 120 mA to 170 mA. (b) QEPAS peak signals as a function of the  $\text{H}_2\text{S}$  concentrations, depicted by black square symbols. The data points were fitted with a linear function, shown by the red line, representing the best linear fit.

$\text{SF}_6$  was continuously flowing in the MMPC. The measured Allan-Werle deviation plot is shown in Fig. 6. Remarkably, a minimum detection limit of  $\sim 25$  ppb can be achieved with an integration time of 100 s. This sensitivity level is achieved in half the time required by a recently reported multi-pass differential photoacoustic sensor, demonstrating the improved performance of our system [44].

For the real-time measurement of  $\text{SF}_6$  decompositions in the GIS system, the response time is a crucial sensor parameter as it determines the accuracy of monitoring and plays a key role in timely maintenance of GIS equipment. A faster response time leads to reduced data delay distortion and enables early detection of GIS faults. The compact design of the enclosed MMPC, with an internal volume of only  $67 \text{ cm}^3$ , facilitates quick exchange of the gas sample, allowing for real-time measurement of  $\text{H}_2\text{S}$  concentration. To evaluate the sensor's response time, a gas flow rate of 70 sccm was chosen, and initially, pure  $\text{SF}_6$  was purged inside the MMPC. The laser drive current was set at 148 mA to lock the laser wavelength to the peak of the  $\text{H}_2\text{S}$  absorption line. Subsequently, a rapid switching mechanism interrupted the flow of pure  $\text{SF}_6$  and initiated the flow of a 100 ppm  $\text{H}_2\text{S}$  in  $\text{SF}_6$  mixture into the MMPC. The rise time refers to the duration required for the LITES signal to increase from the background level to its steady-state value. Finally, the system was switched back to pure  $\text{SF}_6$  to measure the fall time, which represents the duration needed to reach the background signal level. Fig. 7 depicts that

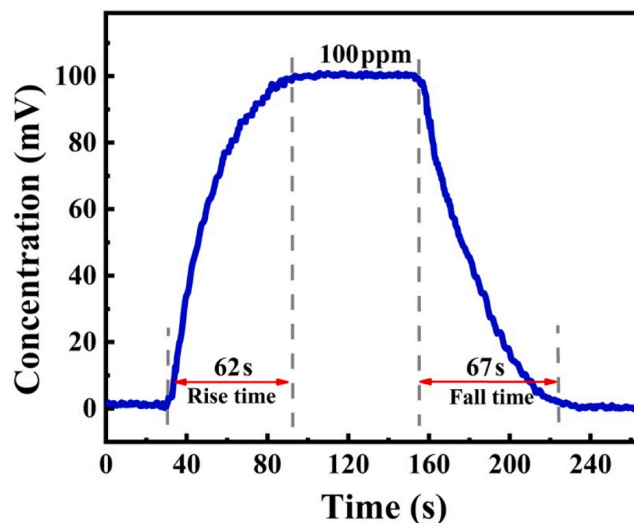


Fig. 7. Sensor response time measurement. Rise and fall times are highlighted in the graph.

the recorded sensor signal during all the measurements. The rise time and fall time were determined to be  $\sim 62$  s and  $\sim 67$  s respectively, which aligns well with the estimated MMPC sample exchange time of  $\sim 57$  s (chamber volume divided by gas flow rate). The test was repeated multiple times without any discernible differences.

## 5. Conclusions

In this work, we reported on a LITES sensor capable of detecting  $\text{H}_2\text{S}$  concentrations in a  $\text{SF}_6$  matrix. LITES, as a non-contact sensing technique, is well-suited for applications that involve the analysis of corrosive gas samples, such as decomposed  $\text{SF}_6$  matrices. The sensor setup consisted of a  $1.58 \mu\text{m}$  DFB laser as the excitation source and a standard QTF coupled with a compact multi-pass cell as the photodetector. The sensor exhibited a linear dynamic range of 0–100 ppm, with an achieved MDL of  $\sim 300$  ppb and an integration time of 300 ms. This sensitivity improved to  $\sim 25$  ppb for a longer integration time of 100 s. The sensor demonstrated a response time of  $\sim 1$  min for a gas flow rate of 70 sccm. These achieved performance metrics meet the requirements for monitoring GIS systems. In future developments, erbium-doped fiber amplifiers could be implemented to enhance the laser power and further improve the detection sensitivity.

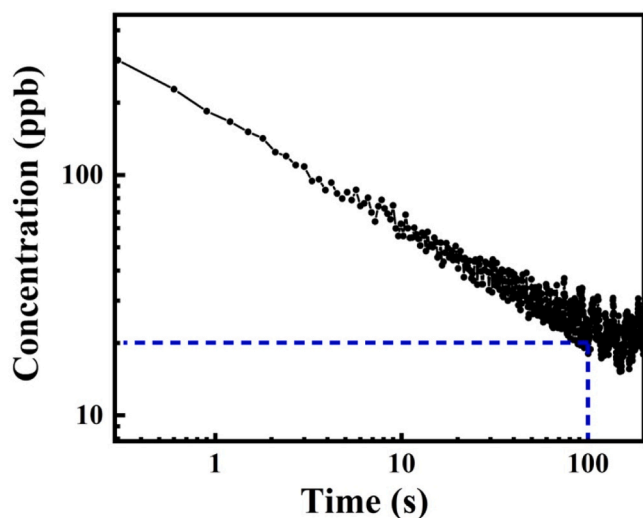


Fig. 6. Allan-Werle deviation plot of the QEPAS signal measured in pure  $\text{SF}_6$  as a function of the integration time. For a 100 s integration time, the minimum detection limit can be significantly reduced to  $\sim 25$  ppb.

## Declaration of Competing Interest

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

## Data Availability

Data will be made available on request.

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