



Contents lists available at ScienceDirect

## Sensors and Actuators: B. Chemical

journal homepage: [www.elsevier.com/locate/snb](http://www.elsevier.com/locate/snb)

## Overcoming dynamic-range limitations in dual-gas sensing via dual-mode QTF excitation

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## ARTICLE INFO

## Keywords:

Dual-spectral  
Multi-gas detection  
LITES  
QEPAS

## ABSTRACT

Simultaneous detection of multiple gases with concentrations spanning five to six orders of magnitude remains a significant challenge for conventional sensing technologies. In this work, we present a dual-spectral (QEPAS-LITES) multi-gas detection system based on a single customized quartz tuning fork (QTF). By applying second-harmonic ( $2f$ ) wavelength modulation at 1368.6 nm and 1653.7 nm, the fundamental and first-overtone resonance modes of the QTF are excited concurrently, enabling intrinsic frequency-domain separation and parallel demodulation of photoacoustic and photothermal signals. This streamlined architecture enables truly simultaneous monitoring of percent-level  $\text{H}_2\text{O}$  and ppm-level  $\text{CH}_4$ , effectively overcoming the dynamic-range limitations of standard multi-gas sensing. Excellent linearity is obtained for both species ( $R^2 = 0.9999$  for  $\text{H}_2\text{O}$  and 0.998 for  $\text{CH}_4$ ), with minimum detection limits of approximately 30 ppm and 660 ppb, respectively. Allan-Werle deviation analysis confirms stable long-term operation for both detection channels. By overcoming the dynamic-range and saturation limitations of conventional multi-gas sensors, the proposed dual-mode approach provides a robust and scalable platform for wide-dynamic-range gas sensing in applications including atmospheric greenhouse gas monitoring, industrial emission analysis, and complex open-path environments.

## 1. Introduction

Online multi-component gas detection plays a crucial role in a wide range of fields, including environmental greenhouse gas monitoring, industrial process safety, and biomedical breath analysis [1–5]. A major challenge arises when target species coexist at concentrations differing vastly different concentrations, often spanning five to six orders of magnitude. For example, atmospheric monitoring requires detecting trace methane ( $\text{CH}_4$ , at  $\sim 1.9$  ppm) in the presence of abundant water vapor ( $\text{H}_2\text{O}$ , at 1%–5%), while industrial emission analysis may involve simultaneous measurement of ppb-level pollutants (e.g.,  $\text{NO}_x$ ) alongside percent-level  $\text{CO}_2$ . Such scenarios impose stringent requirements for sensing technologies: high sensitivity for trace components, strong resistance to saturation for dominant species, robust, interference-free discrimination of multi-species signals. Nevertheless, conventional gas-sensing techniques face fundamental limitations in meeting these

demands. Electrochemical sensors, though compact and low-cost, exhibit poor selectivity, substantial cross-interference (e.g.,  $\text{H}_2\text{O}$  affecting  $\text{CH}_4$  readings), and limited operational lifetimes [6–7]. Semiconductor sensors offer fast response times but are susceptible to environmental fluctuations and lack molecular specificity [8–9]. Among optical methods, tunable diode laser absorption spectroscopy (TDLAS) provides excellent selectivity, yet their dynamic range is constrained ( $\sim 3$ – $4$  orders of magnitude) and signal saturate at high concentrations [10–13]. Cavity-enhanced methods, including cavity ring-down spectroscopy (CRDS) [14] and cavity-enhanced absorption spectroscopy (CEAS) [15] achieve outstanding sensitivity (down to ppb-level) but rely on enclosed high-finesse cavities, which hinder sensing applications and are prone to mirror contamination. Conventional photoacoustic spectroscopy (PAS) [16–21] offers strong immunity to optical interference but depends on bulky acoustic microphones and direct gas-sensor contact, rendering it unsuitable for corrosive or open-path environments.

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Received 3 February 2026; Received in revised form 3 April 2026; Accepted 16 April 2026

Available online 17 April 2026

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Quartz-enhanced photoacoustic spectroscopy (QEPAS) overcomes several limitations of conventional photoacoustic methods by employing a miniaturized quartz tuning fork (QTF, typically  $\sim 5 \times 3 \times 0.25 \text{ mm}^3$ ) as a high-Q acoustic detector, enabling compact, robust, and selective gas detection across multiple species [22–30]. However, standard QEPAS architectures are primarily designed for gas mixtures with relatively narrow concentration ranges (ppm-percent). Their performance degrades substantially in scenarios where trace gases coexist with highly abundant components. Strong photoacoustic signals generated by dominant species can saturate the QTF response, masking weak signals from trace analytes and compromising quantitative accuracy. Additionally, cross-sensitivity and spectral overlap hinder reliable multi-species discrimination under high dynamic-range conditions. A further limitation arises from the contact-based nature of QEPAS direct exposure of the QTF to corrosive gases such as  $\text{NO}_x$ ,  $\text{H}_2\text{S}$  or  $\text{SO}_2$  can lead to electrode degradation, frequency drift, reduced quality factor, and ultimately sensor failure [31–32]. The intrinsic need for physical interaction with the gas sample restricts the use of QEPAS in applications requiring remote sensing capability or operation in chemically aggressive or strongly concentration-stratified environments. To mitigate these drawbacks, laser-induced thermoelastic spectroscopy (LITES) was introduced [33–36]. Based on the detection of laser-induced thermoelastic waves in solids, LITES eliminates the need for mechanical resonators and enables rapid response, high mechanical robustness, and inherent remote-sensing capability. However, its sensitivity to trace gases is generally lower than that of QEPAS due to inherent thermoelastic signal attenuation.

To leverage the high sensitivity of QEPAS together with the sensing and robustness advantages of LITES, several hybrid configurations have been investigated. In 2021, Hu et al. [37] proposed a dual-QTF configuration (QEPAS-PTS) capable of generating both photoacoustic and photothermal signals; however, the system remained restricted to single-gas detection and suffered from intrinsic frequency mismatch between the two tuning forks. Qiao et al. [38] later introduced a single-QTF dual-signal (S-QEDS) scheme that eliminated the mismatch issue, but its functionality was still limited to one target species. In 2023, Cui et al. [39] presented a multi-pass cell (MPC) enhanced dual-spectroscopy method that increased the  $\text{C}_2\text{H}_2$  signal by nearly 80-fold, yet the persistent problem of crosstalk between photoacoustic and photothermal channels remained unsolved. Most recently, Liang et al. [40] developed a QEPAS-LITES hybrid sensor for dissolved-gas monitoring in deep-sea environments. Although their system improved stability for  $\text{CH}_4$ ,  $\text{CO}_2$ ,  $\text{H}_2\text{S}$ , and  $\text{NH}_3$  by using LITES as an auxiliary correction for QEPAS, it neither enabled true simultaneous multi-species detection nor supported open-path operation, as it relied on closed measurement cell.

In this work, we present an integrated dual-mode QEPAS-LITES platform utilizing a single customized quartz tuning fork (QTF) to overcome the dynamic-range limitations of conventional sensors. By concurrently exciting the fundamental and first-overtone resonance modes of the QTF, the system achieves intrinsic frequency-domain isolation, enabling the parallel, crosstalk-free demodulation of photoacoustic and photothermal signals. This unique architecture effectively bypasses signal saturation in high contrast gas mixtures, allowing for the truly simultaneous detection of percent-level  $\text{H}_2\text{O}$  and ppm-level  $\text{CH}_4$ . Compared to multi-sensor configurations, our approach significantly streamlines the optical layout and calibration complexity. We demonstrate a robust, scalable platform that provides a superior solution for wide dynamic range, multi-species monitoring in complex and concentration stratified gas environments.

## 2. QTF Fundamental mode and overtone mode theoretical background

To ensure effective excitation and detection of gas signals, the modulation frequency of the optical excitation must remain much lower

than the molecular relaxation rate of the target species. This requirement imposes a constraint on the acoustic sensor: the resonance frequency of the QTF must satisfy  $f \ll 1/(2\pi\tau)$ , where  $\tau$  is the relaxation time. In practice, QTF operating above  $\sim 50 \text{ kHz}$  fail to meet this condition for many slow-relaxing gas molecules. Standard commercial quartz tuning forks, however, have a fundamental resonance frequency of approximately 32.76 kHz, a quality factor (Q value) of around 10000, while their first overtone mode lies above 190 kHz [41]. Consequently, the high-order modes of conventional QTFs are unsuitable for QEPAS operation because they violate the relaxation-frequency constraint and cannot be efficiently driven by wavelength-modulated photoacoustic excitation. To overcome this limitation, we employed a customized QTF specifically designed with substantially lower resonance frequencies. This device exhibits fundamental and overtone modes in the few-kHz to tens-of-kHz range [42], characterized by its larger dimensions and lower resonance frequencies compared to standard QTFs, was intentionally selected to optimize the excitation of both the fundamental and the first-overtone modes. This design ensures that both modes possess high Q-factors and a wide frequency spacing, providing the necessary frequency-domain orthogonality required for the simultaneous, crosstalk-free detection of gases with vastly different concentration levels. Structurally, the custom QTF is approximately 4.6 times larger than a standard 32-kHz fork and features a high aspect ratio of 17:1. Specifically, this custom design has a prong length ( $L$ ) of 17 mm and a prong thickness ( $T$ ) of 1 mm. Additionally, it features a crystal width ( $w$ ) of 0.25 mm and a prong spacing of 0.7 mm. The resonance frequency  $f_n$  of the  $n$ th vibration mode of the QTF can be expressed as

$$f_n = K \times \frac{T}{L^2} \times n^2$$

where  $K$  is a material-dependent constant (determined by quartz density and Young's modulus),  $T$  is the prong-tip width, and  $L$  is the prong length. Numerical analysis shows that  $n = 1.194$  corresponds to the fundamental mode, whereas  $n = 2.988$  represents the first overtone mode of the customized fork.

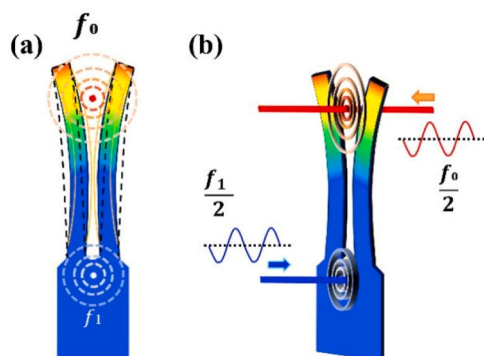
Previous studies [24–25] successfully simulated both the fundamental and the 1st overtone vibration modes of the customized quartz tuning fork, demonstrating that a multi-gas sensor based on frequency-division multiplexing can simultaneously exploit these two resonances to detect different gas species with high stability and sensitivity. Building on this foundation, we propose a dual-mode approach that assigns the fundamental resonance to photoacoustic detection and the first overtone to photothermal detection. In the proposed configuration, one modulated laser beam is focused in the gap between the tuning fork prongs near to their top to generate a photoacoustic response. A second modulated laser beam is directed onto the uncoated quartz crystal region at the base of the tuning fork vibrating arms, inducing a thermoelastic photothermal response.

Fig. 1 illustrates a possible scheme for the proposed dual-mode detection scheme. In Fig. 1(a), the simulated vibrational pattern shows the simultaneous excitation of the fundamental and the 1st overtone modes, which are responsible for generating the photoacoustic and photothermal signals, respectively. Fig. 1(b) depicts the combined vibration state produced when two independently modulated laser beams are incident at their designated excitation positions, depicting how both modes can be concurrently driven without mutual interference.

## 3. Description of the sensor system

### 3.1. Schematic diagram of the sensor

Fig. 2 presents the main modules of the dual-spectral multi-gas detection system, which utilizes a single customized quartz tuning fork (QTF) with a fork 17-mm prong length as the unified signal receiver. The setup consists of two functional subsystems, i.e., a QEPAS and a LITES



**Fig. 1.** (a) Simulated displacement field showing the combined excitation of the fundamental and first-overtone vibration modes; (b) Simulated response of the QTF when two independently modulated laser beams are simultaneously incident at the designated excitation points, illustrating the generation of the combined vibration state.

modules. In the QEPAS section, a DFB laser (Laser 1, L1) mounted on a temperature- and current-controlled driver board (Driver 1, D1) delivers the excitation beam. After collimation, the beam is directed between the QTF prongs and excites either the fundamental or first-overtone mode to generate the corresponding photoacoustic signal. The customized QTF is housed inside a gas chamber equipped with two  $\text{CaF}_2$  windows for optical access. Despite minor Fresnel reflection losses at the interfaces, these windows maintain a high transmittance of over 93.5% in the near-infrared region, guaranteeing sufficient optical power delivery with negligible impact on measurement accuracy. The chamber inlet is connected to a controlled  $\text{H}_2\text{O}:\text{N}_2$  gas mixture, regulated via a mass-flow controller (Alicat M-500SCCM-D) and a needle valve. At the outlet, a vacuum pump (KNF Technologies Ltd., N816.3KT.18) maintains the desired flow rate and chamber pressure, ensuring stable operating conditions.

For the LITES module, a second DFB laser (Laser 2, L2), controlled by a stabilized driver (Driver 2, D2), serves as the excitation source. The laser beam passes through a modular gas interface, currently configured with a 15-cm absorption cell (capable of integrating a multipass cell or other cell types for scalable sensitivity). The beam is then focused by a 50-mm focal-length lens onto the quartz tuning fork surface. Remarkably, the entire detection system fits within a volume of  $\sim 25 \times 6 \times 8 \text{ cm}^3$ . This highly integrated design drastically reduces the sensor's physical footprint relative to conventional QEPAS/LITES setups, facilitating deployment in space-constrained environments.

Fig. 3 illustrates the selected lasers focusing positions for photoacoustic and photothermal excitations. Previous experiments have shown that the strongest QEPAS response for the fundamental mode occurs when the laser beam is focused near the upper region of the QTF [42]. The same experiments demonstrate that the overtone mode shows two antinode positions, one located 9.5 mm from the QTF top and

another one located near the QTF top. Since there is no big difference in terms of signal intensity between these two antinode positions, to avoid adjustments of the laser alignment when changing detection channels, we prefer to fix the laser focusing position for QEPAS (L1) at 2.0 mm from the prong tip.

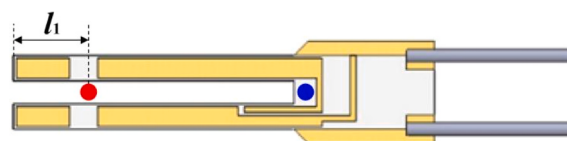
Thereby, 1368.6 nm beam (L1) has been focused between the QTF prongs near the upper region of the fork, approximately 2.0 mm from the prong tip, to optimally excite the fundamental or the 1st overtone mode for  $\text{H}_2\text{O}$  sensing. In contrast, the 1653.74 nm beam (L2) is focused onto the base of the tuning fork, specifically on the uncoated quartz surface at the root of the vibrating arm and under an optimized incidence angle in order to maximize the thermoelastic response [44]. The spatial separation of excitation points, together with the inherent frequency-domain isolation of the two modes, should ensure that the signals remain completely independent and unaffected by each other during simultaneous operation.

A commercial gas dilution instrument (MCQ instruments GB100 series) supplies  $\text{CH}_4:\text{N}_2$  mixtures of controlled concentration to the absorption cell. The inlet flow is regulated using a mass-flow controller (Alicat M-500SCCM-D) and a needle valve, while a vacuum pump (KNF N816.3KT.18) at the outlet maintains stable pressure and flow conditions.

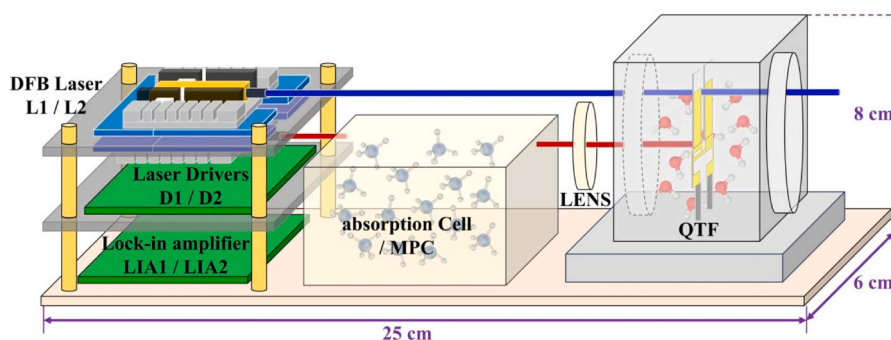
The customized QTF is interfaced with a low-noise trans-impedance amplifier (TA), and the amplified signal is routed to a lock-in amplifier board for second-harmonic ( $2f$ ) demodulation, with demodulation frequencies set to the fundamental and first-overtone resonance modes, respectively. Both detection channels employed a 12 dB/octave filter slope and all the experiments reported in this work have been performed by fixing the integration time to 1 s. The processed  $2f$ , corresponding to photoacoustic and photothermal signals, are recorded and displayed via the data acquisition interface.

### 3.2. Gas absorption characteristics at 1368.6 nm and 1653.74 nm

To evaluate the performance of the proposed dual-mode system under challenging real-world conditions, atmospheric water vapor ( $\text{H}_2\text{O}$ ) and methane ( $\text{CH}_4$ ) were selected as the target gases. This specific gas pair was deliberately chosen to simulate an "extreme dynamic range" scenario typical of environmental monitoring.  $\text{H}_2\text{O}$  is a ubiquitous background gas present at massive, percent-level concentrations (typically 1%–5%), whereas  $\text{CH}_4$  is a slow-relaxing greenhouse gas found at



**Fig. 3.** Laser incidence position to excite the fundamental vibrational mode (red dot) and the 1st overtone vibrational mode (blue dot),  $l_1 = 2.0 \text{ mm}$ .



**Fig. 2.** Main detection modules of the dual spectral multi-gas simultaneous detection system.

trace ppm levels. Additionally, in photoacoustic spectroscopy, H<sub>2</sub>O acts as a strong vibrational-translational (V-T) relaxation promoter for CH<sub>4</sub>. This physical dependence means that even if the CH<sub>4</sub> concentration remains strictly constant, any fluctuation in ambient water vapor will significantly alter the CH<sub>4</sub> photoacoustic signal amplitude, leading to measurement distortions. Therefore, H<sub>2</sub>O must be monitored simultaneously to accurately calibrate the CH<sub>4</sub> measurements in real-world environments. Detecting a weak trace target in the presence of an overwhelming and interfering background is a major bottleneck for conventional sensors. This complex interplay makes this specific gas pair ideal for rigorously validating the system's capacity for crosstalk-free, simultaneous detection across several orders of magnitude. For the experimental implementation, Laser 1 was dedicated to H<sub>2</sub>O detection via the photoacoustic excitation. By stabilizing its temperature and current using DL1, the laser was precisely tuned to 1368.6 nm, delivering an output power of 8.8 mW at an operating temperature of 23 °C and an injected current of 94.5 mA. This emission wavelength is resonant with a strong H<sub>2</sub>O absorption peak at 7306.75 cm<sup>-1</sup>. Although CH<sub>4</sub> exhibits weak transitions within this spectral region, their line strengths are negligible relative to those of H<sub>2</sub>O, ensuring minimal interference from CH<sub>4</sub> during H<sub>2</sub>O measurements [43]. Laser 2 dedicated to the photothermal excitation, was tuned to 1653.74 nm using Driver 2 (DL2), providing an output power of 13.8 mW at 32 °C and an injected current of 81.1 mA. In this spectral region, CH<sub>4</sub> exhibits strong absorption features, whereas H<sub>2</sub>O absorption lines are significantly weaker and do not compromise the accuracy of CH<sub>4</sub> detection. This deliberate selection of two well-separated absorption bands, coupled with precise wavelength stabilization, effectively suppresses mutual spectral interference and enables reliable, mode-specific detection of H<sub>2</sub>O and CH<sub>4</sub>.

## 4. Data analysis of the sensor

### 4.1. Independence between the two different vibration mode

To verify the independence of the fundamental and first-overtone modes, we first measured the resonance response of the customized QTF under atmospheric pressure. A 150 mV sinusoidal excitation was applied to one electrode of the QTF, while the resulting piezoelectric signal was collected from the opposite electrode. By sweeping the excitation frequency and recording the corresponding amplitude response, the resonance peaks were identified. As shown in Fig. 4, the fundamental resonance frequency was measured to be  $f_0 = 2883.8$  Hz, while the first-overtone resonance occurred at  $f_1 = 17839.5$  Hz.

Secondly, we experimentally verified the independence of fundamental frequency and 1st overtone mode, as well as the feasibility of simultaneously demodulating photothermal and photoacoustic signals. 5% CH<sub>4</sub> gas mixture was injected into the gas absorption cell and the pressure was stabilized at 686 Torr. At the same time, L2 was modulated

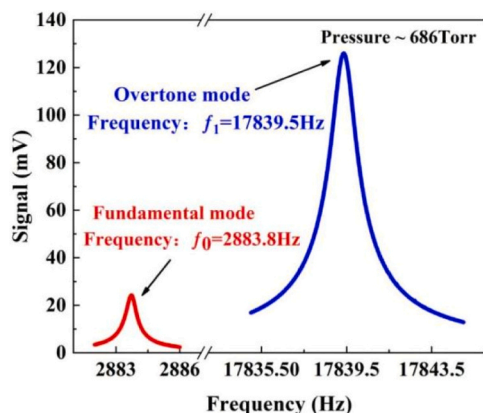


Fig. 4. Resonance spectra of fundamental and 1st overtone mode at 686 Torr.

at half of the fundamental resonance frequency ( $f_0/2$ ) with an 8 mA modulation depth and a 600 s scan period. The resulting photothermal response at the QTF's overtone frequency was extracted using lock-in amplifier LIA 2. Similarly, 1.6% H<sub>2</sub>O was injected into the QTF chamber at the same pressure (686 Torr). L1 was driven with a low-frequency (20 mHz) ramp voltage superimposed on a sinusoidal modulation at  $f_1/2$ , with a 2 mA modulation to acquire the 1st overtone wavelength-modulated photoacoustic signal. To demonstrate the capability for simultaneous, interference-free operation, both lasers were then activated concurrently, each at its designated modulation frequency. As shown in Fig. 5(a-b), the QTF produced a clean photothermal signal of 2.5 mV for 5% CH<sub>4</sub> at the fundamental mode and a distinct photoacoustic signal of ~1.64 mV for 1.6% H<sub>2</sub>O at the 1st overtone mode. We also inverted the excitation scheme, using the fundamental mode for photoacoustic excitation and the first-overtone mode for photothermal excitation. Fig. 5(c-d) shows the simultaneous signals measured for H<sub>2</sub>O (peaks of ~830  $\mu$ V) and CH<sub>4</sub> (peak of ~5.6 mV). In all operating conditions, the CH<sub>4</sub> spectrum remains unchanged when L2 is switched off, and likewise the H<sub>2</sub>O spectrum is unaffected when L1 is disabled. This behavior confirms the complete spectral and modal isolation of the two detection channels.

Based on these observations, we adopted a mode-selective strategy for the subsequent experiments: the first-overtone mode combined with the photothermal effect was used for CH<sub>4</sub> detection, while the fundamental mode combined with the photoacoustic effect was exploited for H<sub>2</sub>O detection.

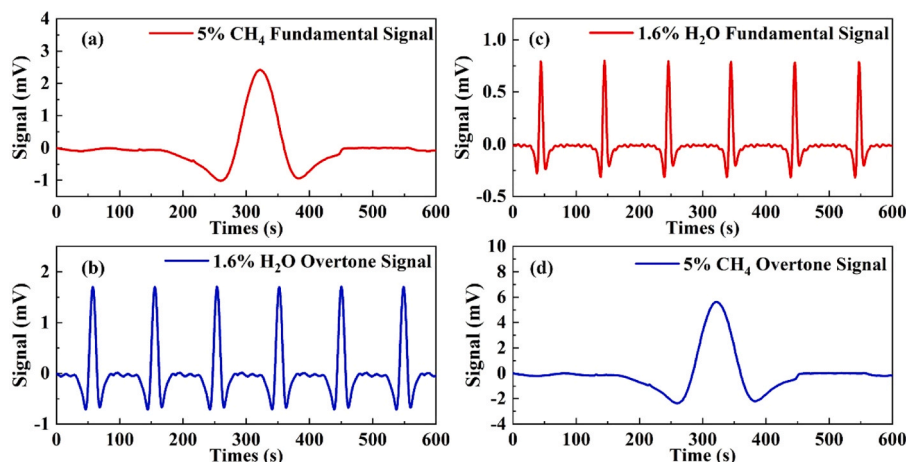
### 4.2. Optimization of gas pressure

Previous studies have shown that the QEPAS signal exhibits a pressure-dependent maximum because decreasing pressure increases the QTF quality factor while reducing acoustic generation efficiency, leading to an optimal trade-off that maximizes signal amplitude [42]. Likewise, the sensitivity of LITES depends on the pressure of the gas mixture surrounding the QTF. As the pressure decreases, air damping is reduced and the Q-factor of the tuning fork increases, thereby increasing the oscillator accumulation time  $\tau = Q/(2\pi f)$ . A longer accumulation time enhances the thermoelastic response, resulting in a marked increase in the LITES signal-to-noise ratio at lower pressures [44].

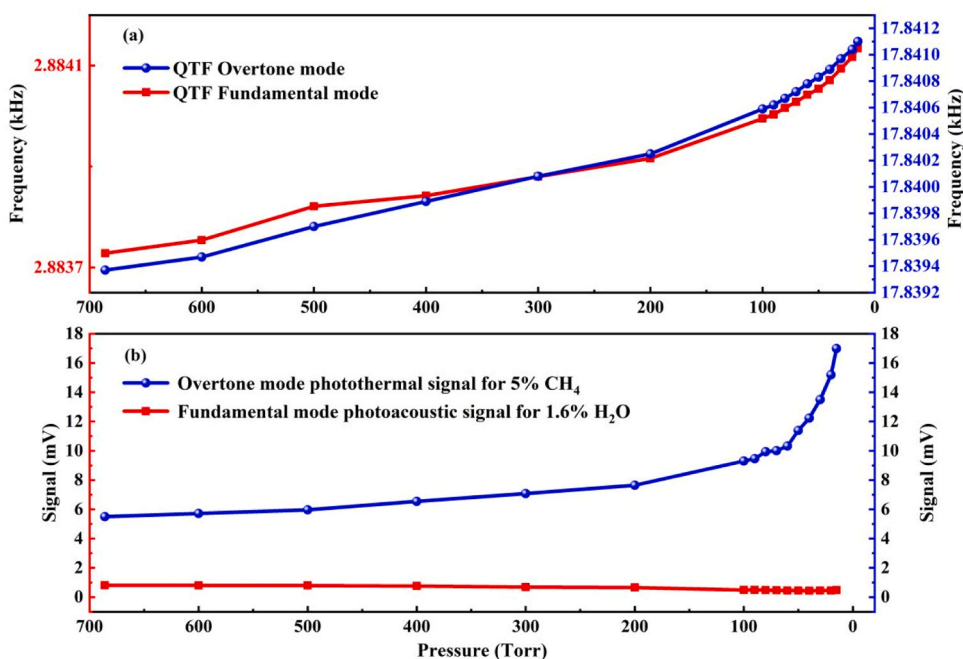
To identify the optimal operating pressure for simultaneous detection, we varied the QTF chamber pressure from 686 Torr to 15 Torr and measured the QTF resonance response. As shown in Fig. 6(a), both the fundamental and first-overtone resonance frequencies increased as the pressure decreased, with a particularly pronounced rise below 100 Torr. At the same time, the Q-factor exhibited a substantial improvement. For example, at 15 Torr the fundamental mode the Q-factor rises up to 15,270, about 2.5 times higher than its value at 686 Torr, while for the first-overtone mode the Q-factor reach a value of 32,767, representing a ~2.15-fold increase. These trends are expected since reduced pressure enhances QTF resonance sharpness.

Subsequently, simultaneous measurements for a gas mixture of 5% CH<sub>4</sub>:N<sub>2</sub> in the absorption cell at a pressure of 686 Torr, while mixtures of 1.6% H<sub>2</sub>O:N<sub>2</sub> filled at different pressures the QTF chamber. As shows in Fig. 6(b), the CH<sub>4</sub> photothermal signals increased steadily as the pressure in the QTF chamber decreases, rising from 5.6 mV to ~17.00 mV. In contrast, the H<sub>2</sub>O photoacoustic signal exhibited a decreasing trend, dropping from ~830  $\mu$ V to ~450  $\mu$ V as pressure inside the QTF chamber was reduced. These opposite pressure-dependent behaviors are fully consistent with previous observations, and demonstrated that the system can operate to simultaneously detect two gases with widely different concentrations using a single customized QTF.

To investigate the sensor system performance, additional measurements were conducted at a fixed pressure of 15 Torr in the QTF chamber. Different CH<sub>4</sub>:N<sub>2</sub> mixtures with methane concentrations varying from 50 ppm to 500 ppm were then flushed into the absorption cell, while keeping fixed the pressure to 686 Torr, and the CH<sub>4</sub> photothermal



**Fig. 5.** Independent and simultaneous detection of photoacoustic and photothermal signals using the fundamental and first-overtone modes of the customized QTF. (a) Fundamental-mode photothermal signal generated by 5% CH<sub>4</sub>; (b) First-overtone-mode photoacoustic response to 1.6% H<sub>2</sub>O; (c) Fundamental-mode photoacoustic signal generated by 1.6% H<sub>2</sub>O; (d) First-overtone-mode photothermal response to 5% CH<sub>4</sub>.



**Fig. 6.** (a) Resonance frequencies of QTF fundamental and 1st overtone vibrational modes measured at different pressure; (b) Photothermal signal measured for 5% CH<sub>4</sub>:N<sub>2</sub> with the QTF operating at the 1st overtone mode and photoacoustic signal of 1.6% H<sub>2</sub>O:N<sub>2</sub> while operating the QTF at the fundamental mode, as a function of the surrounding gas pressure.

response was recorded operating the QTF at its 1st overtone mode. The measured signal increases linearly with CH<sub>4</sub> concentration from  $\sim 0.2$  mV (for 50 ppm) up to  $\sim 1.64$  mV (for 500 ppm). Linear regression of the concentration-dependent data yielded an  $R^2$  value of 0.998 (Fig. 7 (a)), demonstrating excellent linearity of the CH<sub>4</sub> photothermal response in the investigated methane concentration range.

During the CH<sub>4</sub> detection, the H<sub>2</sub>O photoacoustic signal was simultaneously recorded using the fundamental mode. Gas mixture with H<sub>2</sub>O concentrations of 0.4%, 0.6%, 0.8%, 1.0%, 1.2%, and 1.6% were introduced into the QTF chamber, keeping fixed the pressure at 15 Torr. The signal increased linearly with water vapor concentration ( $R^2 = 0.9999$ ), reaching a value of  $\sim 450$   $\mu$ V for 1.6% H<sub>2</sub>O. The corresponding calibration curve is shown in Fig. 7(b) (red line). Considering the detection bandwidth of 0.833 Hz, 12 dB filter slope and 1 s integration time, the minimum detection limit were calculated to be  $\sim 660$  ppb for

CH<sub>4</sub> and  $\sim 30$  ppb for H<sub>2</sub>O.

To assess the long-term stability of the dual-spectral system, Allan-Werle deviation analysis was performed for both detection channels for a pure N<sub>2</sub> gas sample (Fig. 7(c)-(d)), while keeping the pressure fixed at 686 Torr in the absorption cell and 15 Torr in the QTF chamber. This continuous monitoring lasted for over 1.5 h, providing clear evidence of the system's long-term stability without any observable mode-coupling or resonance drift. The results confirm that sensor noise is dominated by thermal noise up to the averaging time of  $\sim 80$  s for photothermal and  $\sim 200$  s for photoacoustic measurements, implying that further sensitivity improvements can be achieved by extending the integration time.

Finally, the system's modularity capability was validated by substituting the LITES absorption cell with a multipass cell providing an optical pathlength  $\sim 4.2$  m and detecting a 10 ppm CH<sub>4</sub> signal at a pressure of 686 Torr, while simultaneously acquiring the H<sub>2</sub>O

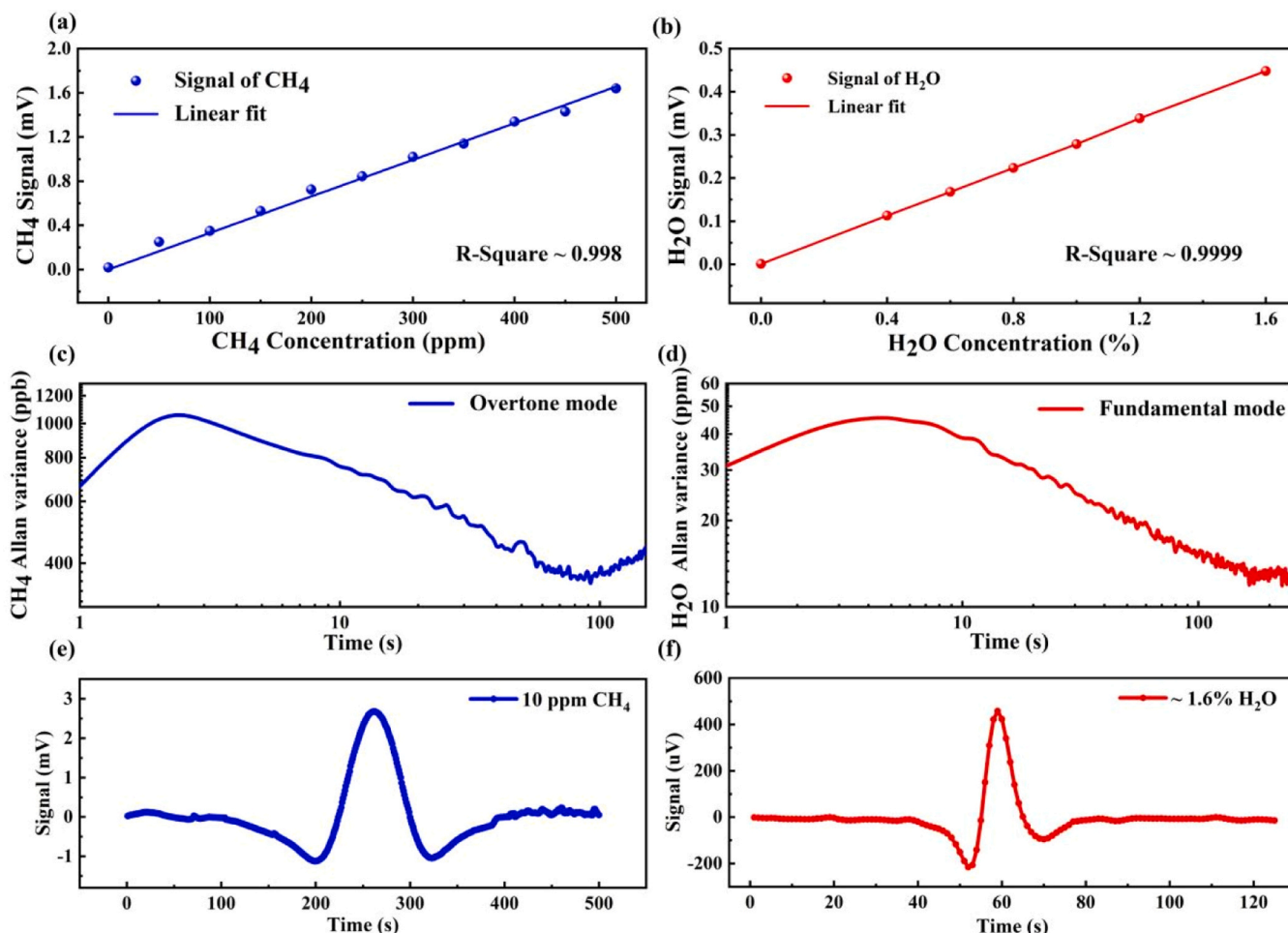


Fig. 7. (a) Photothermal signal (blue dots) measured versus  $\text{CH}_4$  concentration. The blue line represents the best linear fit of the data; (b) photoacoustic signal (red dots) measured versus  $\text{H}_2\text{O}$  concentration. The red line represents the best linear fit of the data; (c) Allan-Werle deviations of the photothermal signal measured for pure  $\text{N}_2$  samples filling the absorption cell at a pressure of 686 Torr, while L2 wavelength was locked to the  $\text{CH}_4$  target absorption line; (d) Allan-Werle deviations of the photoacoustic signal measured for pure  $\text{N}_2$  samples filling the QTF chamber at a pressure of 15 Torr, while L2 wavelength was locked to the target  $\text{H}_2\text{O}$  absorption line; (e)  $2f$  photothermal signal measured for 10 ppm  $\text{CH}_4:\text{N}_2$  at 686 Torr when exploiting a multipass cell with an optical pathlength of  $\sim 4.2$  m; (f)  $2f$  photoacoustic signal simultaneously measured for 1.6%  $\text{H}_2\text{O}:\text{N}_2$  at 15 Torr filling the QTF chamber.

photoacoustic signal. The measured spectra are shown in Fig. 7(e)–(f), with a peak value measured for methane of  $\sim 2.75$  mV, demonstrating the modularity of the proposed simultaneous dual-gas detection system.

## 5. Conclusions

In this work, we demonstrated a dual-spectral multi-gas detection system that combines QEPAS and LITES within a single customized quartz tuning fork. By simultaneously exciting the fundamental and first-overtone resonance modes, the proposed architecture enables intrinsic frequency-domain separation of photoacoustic and photo-thermal responses, allowing truly simultaneous, crosstalk-free detection of gases that differ in concentration by several orders of magnitude. Using  $\text{H}_2\text{O}$  at percent-level and  $\text{CH}_4$  at ppm-level as representative targets, we achieved excellent linearity, high sensitivity, and robust long-term stability. The system also allow to easily exchange the absorption cell type in the LITES module, as demonstrated through 10-ppm methane detection exploiting a multipass cell with an optical pathlength of  $\sim 4.2$  m.

Pressure-dependent studies confirmed that the two modes behave independently and can be optimized separately, enabling efficient signal extraction for both high- and low-concentration gases. The simplified optical configuration, single-resonator design, and strong resistance to spectral and modal interference collectively establish this dual-mode

approach as a versatile and scalable platform for wide-dynamic-range gas sensing.

In future real-world deployment, a sequential flow architecture will be developed. The gas mixture will, for example, first enter the LITES multi-pass absorption cell and subsequently flow into the QEPAS QTF chamber. Since QEPAS and LITES require different optimal operating pressures, a proportional micro-metering valve (or a differential pressure controller) will be implemented in the gas line connecting the two chambers. This configuration will allow employing a single continuous sample intake while actively maintaining the independent optimal pressures.

Overall, the results bridge the performance gap between traditional single-mode QEPAS or LITES systems and the requirements of modern multi-component monitoring in complex or open-path environments. Future developments will include extending the approach to additional molecular species, refining QTF geometries to enhance robustness in harsh environments, and exploring multiplexed operation via higher-order overtone modes. These advancements will further expand the applicability of the technique to challenging scenarios such as industrial emission monitoring, high-temperature process control, and atmospheric sensing.

## CRedit authorship contribution statement

**Hongpeng Wu:** Writing – review & editing, Supervision, Resources, Methodology, Data curation, Conceptualization. **Vincenzo Spagnolo:** Writing – review & editing, Supervision, Methodology, Conceptualization. **Lei Dong:** Writing – review & editing, Supervision, Methodology, Funding acquisition, Conceptualization. **Pietro Patimisco:** Writing – review & editing, Methodology, Conceptualization. **Angelo Sampaolo:** Writing – review & editing, Methodology, Conceptualization. **Xiaowen Shen:** Writing – original draft, Methodology, Investigation, Conceptualization. **Chaofeng Sun:** Methodology, Investigation, Data curation, Conceptualization.

## Declaration of Competing Interest

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

## Acknowledgments

This project was supported by the National Science and Technology Major Project of China (Grant No. 2025ZD1200704); The National Natural Science Foundation of China (NSFC) (Nos. 62475137, 62235010, 62502163, 62501370); Shanxi Provincial Special Fund for Scientific and Technological Cooperation and Exchange (202404041101022, 202304041101019); Fundamental Research Program of Shanxi Province, China (No. 202403021212183, 202303021222034); The Scientific and Technological Innovation Programs of Higher Education Institutions in Shanxi Province of China (2024L014); The National Engineering Research Center of UHV Technology and New Electrical Equipment Basis (Grant No. NERCUIHE-2024-KF-12). Authors from Polysense Lab acknowledge funding from Project PNC 0000001 D3-4-Health—Digital Driven Diagnostics, prognostics and therapeutics for sustainable Health care (CUP: B83C22006120001); MUR – Dipartimenti di Eccellenza 2023–2027 – Quantum Sensing and Modelling for One-Health (QuaSiModO) and THORLABS GmbH within the PolySenSe joint research laboratory.

## Data availability

Data will be made available on request.

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